

5.5.1 Solar wind equ $(V^2 - \frac{2k_b T}{m}) \frac{1}{V} \frac{dV}{dr} = \frac{4k_b T}{m r} - \frac{GM_\odot}{r^2}$

a)

Take given solution $\frac{V^2}{2} - \frac{2k_b T}{m} \ln V = \frac{4k_b T}{m} \ln r + \frac{GM_\odot}{r} + K$

and differentiate wrt r:

$$\frac{2V}{2} \frac{dV}{dr} - \frac{2k_b T}{m} \frac{1}{V} \frac{dV}{dr} = \frac{4k_b T}{m} \frac{1}{r} - \frac{GM_\odot}{r^2}$$

ie $(V^2 - \frac{2k_b T}{m}) \frac{1}{V} \frac{dV}{dr} = \frac{4k_b T}{m} \frac{1}{r} - \frac{GM_\odot}{r^2}$ which is the solar wind equ.

b) At $r = r_c$ $V^2 = \frac{2k_b T}{m} \equiv c_s^2$ so soln is

$$\frac{c_s^2}{2} - c_s^2 \ln c_s - 2c_s^2 \ln r_c - \frac{GM_\odot}{r_c} = K$$

This could be put back into solution to give.

$$\frac{V^2 - c_s^2}{2} - c_s^2 \ln\left(\frac{V}{c_s}\right) = 2c_s^2 \ln\left(\frac{r}{r_c}\right) + GM_\odot \left(\frac{1}{r} - \frac{1}{r_c}\right)$$

c) Density $n = \frac{1}{4\pi m} \frac{1}{V r^2}$ so $\frac{n}{n_\oplus} = \left(\frac{1 \text{ AU}}{r}\right)^2$ since $V = \text{const}$

$$B_r = B_{r\oplus} \left(\frac{1 \text{ AU}}{r}\right)^2 ; B_\phi = -\frac{r}{V} B_{r\oplus} \left(\frac{1 \text{ AU}}{r}\right)^2 = B_{\phi\oplus} \left(\frac{1 \text{ AU}}{r}\right)$$

Planet	r(AU)	n(cm ⁻³)	B _r (nT)	B _{phi} (nT)	B(nT)	psi_spiral(deg)
Earth	1.0000	6.6000	4.9497	-4.9497	7.0000	45.0000
Mercury	0.3900	43.3925	32.5427	-12.6917	34.9300	21.3058
Mars	1.5000	2.9333	2.1999	-3.2998	3.9659	56.3099
Jupiter	5.2000	0.2441	0.1831	-0.9519	0.9693	79.1145
Neptune	30.0000	0.0073	0.0055	-0.1650	0.1651	88.0908

using $B = \sqrt{B_r^2 + B_\phi^2}$ and $\tan \psi_{\text{spiral}} = -\frac{B_\phi}{B_r}$

$S_0 S_0^2$ At standoff $\frac{B^2}{2\mu_0} = \rho V^2 + \frac{B_{imf}^2}{2\mu_0}$

Using $B_{subsolar} = \frac{M}{r^3}$ gives us

$\frac{M^2}{2\mu_0 r_s^6} = \rho V^2 + \frac{B_{imf}^2}{2\mu_0}$ which can be re-arranged

$\frac{M^2}{\mu_0 [\rho V^2 + \frac{B_{imf}^2}{\mu_0}]} = r_s^6$ so $r_s^6 = \frac{M^2}{2\mu_0 \rho V^2 + B_{imf}^2}$

Now $\rho = n m_p = n (\text{cm}^{-3}) 10^6 \frac{\text{kg}}{\text{cm}^3} 1.7 \times 10^{-27} = 1.7 \times 10^{-21} n (\text{cm}^{-3}) \text{ kg/m}^3$

$V^2 = (450 \text{ km/s})^2 = 20 \times 10^{10} = 2 \times 10^{11} (\text{m/s})^2$

$\mu_0 = 4\pi \times 10^{-7}$

$B_{imf}^2 (\text{T}^2) = 10^{-18} B_{imf}^2 (\text{nT})$

Planet	r(AU)	n(cm ⁻³)	Br(nT)	Bphi(nT)	B(nT)	psi_spiral(deg)
Earth	1.0000	6.6000	4.9497	-4.9497	7.0000	45.0000
Mercury	0.3900	43.3925	32.5427	-12.6917	34.9300	21.3058
Mars	1.5000	2.9333	2.1999	-3.2998	3.9659	56.3099
Jupiter	5.2000	0.2441	0.1831	-0.9519	0.9693	79.1145
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Planet	rho(kg/m ³)	Bimf ² (T ²)	M(Tm ³)	rs ⁶ (m)	rs(m)	rplanet(m)	rs/rp
Earth	+1.12E-20	+4.90E-17	+7.90E+15	+1.10E+46	+4.71E+07	+6.37E+06	7.40
Mercury	+7.38E-20	+1.22E-15	+4.60E+12	+5.52E+38	+2.86E+06	+2.49E+06	1.15
Mars	+4.99E-21	+1.57E-17	+1.00E+12	+3.96E+38	+2.71E+06	+3.40E+06	0.80
Jupiter	+4.15E-22	+9.40E-19	+1.50E+20	+1.07E+56	+2.18E+09	+7.14E+07	30.53
Neptune	+1.25E-23	+2.73E-20	+2.00E+17	+6.36E+51	+4.30E+08	+2.48E+07	17.35

Notes:

Mercury just holds off the solar wind

Mars' field is too weak to stand off the solar wind

Jupiter has an enormous magnetosphere (3 R_⊙!)

5.5.3 Re-do derivation starting at (5.16)

$$I = 4\pi r^2 \rho V = \text{const}$$

$$\rho V \frac{dV}{dr} = - \frac{d\rho}{dr} - \frac{\rho GM_{\odot}}{r^2}$$

but $\rho r^{-\gamma} = C$
 $\Rightarrow \rho^{-\gamma} \frac{d\rho}{dr} - \rho \gamma \rho^{-\gamma-1} \frac{d\rho}{dr} = 0$

$$= - \frac{\gamma \rho}{\rho} \frac{d\rho}{dr} - \frac{\rho GM_{\odot}}{r^2} \quad (*)$$

$$\Rightarrow \frac{d\rho}{dr} = + \frac{\gamma \rho}{\rho} \frac{d\rho}{dr}$$

$$\rho = \frac{I}{4\pi r^2 V} \Rightarrow \frac{d\rho}{dr} = \frac{I}{4\pi} \left[-\frac{2}{r^3 V} - \frac{1}{r^2 V^2} \frac{dV}{dr} \right]$$

So momentum eqn becomes

$$\rho V \frac{dV}{dr} = + \frac{\gamma \rho}{\rho} \frac{I}{4\pi} \left[\frac{2}{r^3 V} + \frac{1}{r^2 V^2} \frac{dV}{dr} \right] - \frac{\rho GM_{\odot}}{r^2}$$

$$= \frac{\gamma \rho}{\rho} \cancel{\rho V} \left[\frac{2}{r^3 V} + \frac{1}{r^2 V^2} \frac{dV}{dr} \right] - \frac{\rho GM_{\odot}}{r^2} = \gamma \rho \left[\frac{2}{r} + \frac{1}{V} \frac{dV}{dr} \right] - \frac{\rho GM_{\odot}}{r}$$

Divide by ρ and re-arrange:

$$\left(V^2 - \frac{\gamma \rho}{\rho} \right) \frac{1}{V} \frac{dV}{dr} = \frac{\gamma \rho}{\rho} \frac{2}{r} - \frac{GM_{\odot}}{r^2}$$

Adiabatic Solar wind Equation

To solve this we need to re-write $\frac{\gamma \rho}{\rho} = \gamma C \rho^{\gamma-1} = \cancel{\gamma C \rho^{\gamma-1}}$
 Go back to (*) and divide by ρ to give

$$V \frac{dV}{dr} = - \gamma C \rho^{\gamma-1} \frac{1}{\rho} \frac{d\rho}{dr} - \frac{GM_{\odot}}{r^2} = - \gamma C \rho^{\gamma-2} \frac{d\rho}{dr} - \frac{GM_{\odot}}{r^2}$$

ie $\frac{d}{dr} \left(\frac{1}{2} V^2 \right) = - \gamma C \frac{1}{\gamma-1} \frac{d}{dr} (\rho^{\gamma-1}) + \frac{d}{dr} \left(\frac{GM_{\odot}}{r} \right)$ which can be

integrated to give

$$\frac{V^2}{2} + \frac{\gamma}{\gamma-1} C \rho^{\gamma-1} - \frac{GM_{\odot}}{r} = K$$

$$C \rho^{\gamma-1} = \rho \rho^{-\gamma} \rho^{\gamma-1} = \frac{\rho}{\rho}$$

$$\text{ie } \left[\frac{V^2}{2} + \frac{\gamma}{\gamma-1} \frac{\rho}{\rho} - \frac{GM_{\odot}}{r} = K \right]$$